

**PHYSICS 211**  
**Final EXAM, Fall 2011-2012**  
**TIME: 120 minutes**

January 23, 2012

**DO NOT OPEN THIS EXAM BEFORE YOU ARE TOLD TO BEGIN**

**The usage of programmable calculators is strictly forbidden**

NAME \_\_\_\_\_

ID Number \_\_\_\_\_

Useful information

$$\epsilon_0 = 8.85 \times 10^{-12} \text{ C}^2/\text{N}\cdot\text{m}^2$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ Tm/A}$$

$$q = 1.6 \times 10^{-19} \text{ C}$$

$$m_e = 9.1 \times 10^{-31} \text{ kg}$$

$$m_p = 1.67 \times 10^{-27} \text{ kg}$$

$$k_e = 8.9875 \times 10^9 \text{ Nm}^2/\text{C}^2$$

**Grading**

<b>Part I (20)</b>	
<b>Part II (20)</b>	
<b>Part III (30)</b>	
<b>Part IV (30)</b>	
<b>Total</b>	

Check if solution is continued on the back.

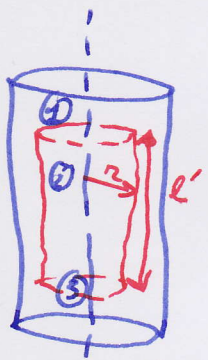
**Part I: Gauss's Law (20)**

Consider a long cylinder with radius  $R$  and a length  $l$  that is very large when compared to  $R$  so as to neglect the cylinder's end edge effects. This long cylinder has a volumetric charge density  $\rho$ .

- (a) (7) Show that the electric field **inside** the cylinder at a distance  $r$  from its axis has the expression  $E = (\rho/2\epsilon_0)r$ .

To determine the electric field, we use Gauss's Law

$$\Phi_E = \int_{G.S.} \vec{E} \cdot d\vec{A} = \frac{Q_{in}}{\epsilon_0}; \text{ The Gaussian surface (G.S.) is a cylinder with radius } r.$$



$$\Phi_E = \Phi_1 + \Phi_2 + \Phi_3 \text{ or } \Phi_1 = \Phi_3 = 0 \text{ because } \vec{E} \perp \vec{A}$$

$$\Phi_2 = E(2\pi r l') = \frac{Q_{in}}{\epsilon_0} \text{ or } \rho \pi r^2 l' = \frac{Q_{in}}{\epsilon_0} \Rightarrow Q_{in} = \rho \frac{\pi R^2 l'}{\pi R^2 l'}$$

$$\Rightarrow \boxed{E = \frac{\rho}{2\epsilon_0} r}$$

$$Q_{in} = \rho \pi r^2 l'$$

- (b) (8) Show that the electric field **outside** the cylinder at a distance  $r$  from its axis has the expression  $E = 2\pi k_e \rho R^2 / r$

The Gaussian surface is now with  $r > R$ , it is a cylinder

$$\Phi_E = \Phi_1 + \Phi_2 + \Phi_3 = E(2\pi r l') = \frac{\rho \pi R^2 l'}{\epsilon_0}$$

$$\text{so } \Phi_1 = \Phi_3 = 0$$

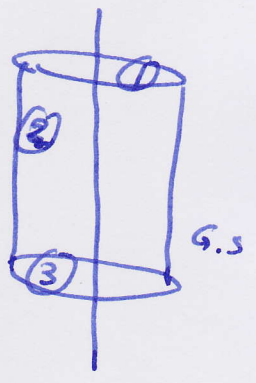
$$\Rightarrow E = \frac{\pi}{2\epsilon_0} \rho \frac{R^2}{r} = 2k_e \rho \frac{R^2}{r}$$

- (c) (5) When  $R \rightarrow 0$  the wire can be described by a linear charge density  $\lambda$ , show that  $\lambda = Q/l$ . [Hint: Compare the electric field above to the one obtained with a linear charge density  $\lambda$ ]

For a thin wire, the Gaussian surface is a cylinder with radius  $r$  and length  $l'$

$$\Phi_E = \Phi_1 + \Phi_2 + \Phi_3; \Phi_1 = \Phi_3 = 0 \text{ as } \vec{E} \perp \vec{A}$$

$$\Rightarrow \Phi_2 = E(2\pi r l') = \frac{Q_{in}}{\epsilon_0} \Rightarrow E = 2k_e \frac{Q_{in}}{r}$$



$$E = 2k_e \frac{Q_{in}}{r} = 2\pi k_e \rho \frac{R^2}{r} \Rightarrow Q_{in} = \pi R^2 \rho = \frac{\pi R^2}{\pi R^2 l} Q = \frac{Q}{l}$$

**Part II: Ampere's Law (20)**

- (a) (4) Assume that the charges of the cylinder in part I are now moving with a drift velocity  $v_d$ , determine the expression of the current  $I$ .

$$I = \frac{dQ}{dt} \text{ or } Q = Nq \text{ with } N = nV = nA\ell$$

$$\Rightarrow I = qnA \frac{dx}{dt} \text{ on the average } \boxed{I = nqAv_d}$$

- (b) (6) Show that the magnetic field **inside** the wire has the expression

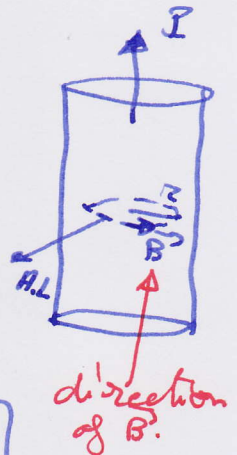
$$B = \left( \frac{\mu_0 I}{2\pi R^2} \right) r ; \text{ make your own drawing to show its direction.}$$

We use Ampere's law:  $\oint_{A.L.} \vec{B} \cdot d\vec{\ell} = \mu_0 I_{in}$   
 We choose the Amperian loop (A.L.) to be a circle with radius  $r$ .

$$\Rightarrow \oint \vec{B} \cdot d\vec{\ell} = \oint B d\ell = B \oint d\ell = B(2\pi r) = \mu_0 I_{in}$$

The current is uniform  $\Rightarrow$

$$j = \frac{I}{A} = \frac{I_{in}}{A_{A.L.}} \Rightarrow I_{in} = \frac{\pi R^2}{\pi r^2} I \Rightarrow \boxed{B = \frac{\mu_0 I}{2\pi R^2} \cdot r}$$



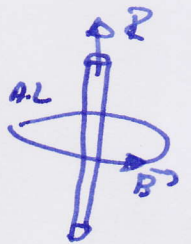
- (c) (5) Show that the magnetic field **outside** the wire has the expression  $B = \frac{\mu_0 I}{2\pi r}$ ;

make a drawing to show its direction. (Same as above)

The Ampere's law:  $\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{in}$

$$\Rightarrow B(2\pi r) = \mu_0 I \text{ as } I_{in} = I$$

$$\Rightarrow \boxed{B = \frac{\mu_0 I}{2\pi r}}$$



- (d) (5) State and verify Gauss's law for magnetism using the wire with current  $I$ .

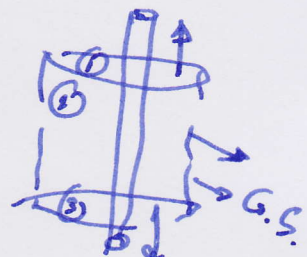
Gauss's law for magnetism reads:  $\Phi_B = \int_{G.S.} \vec{B} \cdot d\vec{A} = 0$

We choose the ~~the~~ Gaussian surface to be a closed cylinder around the wire

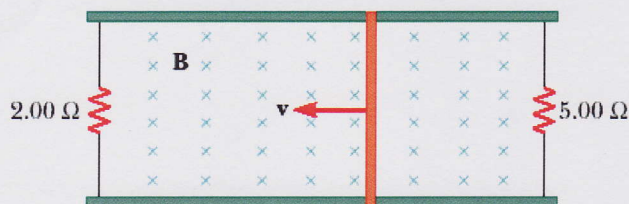
$$\Rightarrow \Phi_B = \Phi_1 + \Phi_2 + \Phi_3$$

$$\Phi_1 = \Phi_2 = \Phi_3 = 0 \text{ as } \vec{B} \perp d\vec{A} \text{ everywhere}$$

$$\Rightarrow \boxed{\Phi_B = 0}$$



A conducting rod of length  $\ell = 35.0$  cm is free to slide on two parallel conducting bars as shown in the Figure below. Two resistors  $R_1 = 2.00 \Omega$  and  $R_2 = 5.00 \Omega$  are connected across the ends of the bars to form a loop. A constant magnetic field  $B = 2.50$  T is directed perpendicularly into the page. An external agent pulls the rod to the left with a constant speed of  $v = 8.00$  m/s.



- (a) **5** Using Faraday's law, find the analytical expression and the value of the emf caused by the motion of the bar.

Faraday's law states  $\epsilon = -\frac{d\Phi_B}{dt}$  where  $\Phi_B$  is the magnetic flux  $\Phi_B = \int \vec{B} \cdot d\vec{A}$ . Take  $\mathbf{A}$  to be a rectangle of length  $l$  and width  $dx$  and taking advantage of the fact that the magnetic field is constant leads to:

$$\epsilon = -\frac{d\Phi_B}{dt} = -Bl \frac{dx}{dt} = -Blv$$

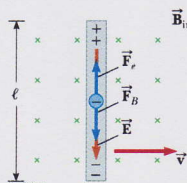
$\epsilon = Blv = (2.50 \text{ T})(0.350 \text{ m})(8.00 \text{ m/s}) = 7.00 \text{ V}$

- (b) **6** Consider the motion of free electrons in the bar. Show that the same expression of the emf can be found when equilibrium is reached between the electric and the magnetic forces.

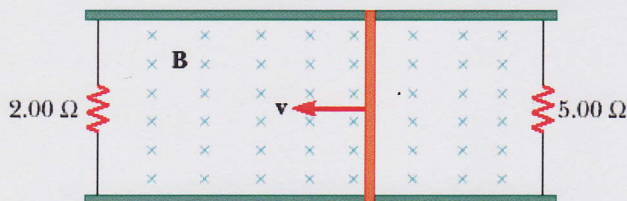
Two forces are acting on free electrons, the magnetic force  $\vec{F} = q\vec{v} \times \vec{B}$  and the electric force. At equilibrium, we have the sum of the forces is equal to 0 leading to  $qE = qvB$  or  $E = vB$

The electric field is related to the potential difference across the ends of the conductor:

$$\Delta V = E\ell = B\ell v$$



- (c) **4** Show on the figure above the direction of the currents in the two resistors, and justify below.



The left-hand loop contains decreasing flux away from you, so the induced current in it will be clockwise, to produce its own field directed away from you. Let  $I_1$  represent the current flowing

upward through the  $2.00\text{-}\Omega$  resistor. The right-hand loop will carry counterclockwise current. Let  $I_3$  be the upward current in the  $5.00\text{-}\Omega$  resistor.

- (d) **4** Determine the value of the current in both resistors.

Kirchhoff's loop rule then gives:

$$+7.00\text{ V} - I_1(2.00\ \Omega) = 0$$

$$I_1 = \boxed{3.50\text{ A}}$$

and  $+7.00\text{ V} - I_3(5.00\ \Omega) = 0$   $I_3 = \boxed{1.40\text{ A}}$ .

- (e) **5** Find the total power delivered to the resistance of the circuit.


The total power dissipated in the resistors of the circuit is

$$P = \varepsilon I_1 + \varepsilon I_3 = \varepsilon(I_1 + I_3) = (7.00\text{ V})(3.50\text{ A} + 1.40\text{ A}) = \boxed{34.3\text{ W}}$$

- (f) **6** Find the magnitude of the applied force that is needed to move the rod with this constant velocity.

*Method 1:* The current in the sliding conductor is downward with value

$$I_2 = 3.50\text{ A} + 1.40\text{ A} = 4.90\text{ A}. \text{ The magnetic field exerts a force of}$$

$$F_m = \ell B = (4.90\text{ A})(0.350\text{ m})(2.50\text{ T}) = 4.29\text{ N}$$
 directed  toward the right on this conductor. An outside agent must then exert a force of

$$\boxed{4.29\text{ N}}$$
 to the left to keep the bar moving.

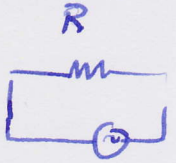
**Part IV: AC circuits (30)**

We consider a circuit composed of an AC power source, supplying a current  $I = I_{\max} \sin(\omega t)$ , a resistor  $R$  and an inductor  $L$ .

(a)(4) Consider that the **resistance alone** is connected to the power supply, determine the expression of the potential difference across it.

We use Ohm's law  $\Delta V = RI$

$$\Rightarrow \Delta V(t) = R I_{\max} \sin \omega t \Rightarrow \begin{cases} \Delta V_{\max} = R I_{\max} \\ \phi = 0 \end{cases}$$

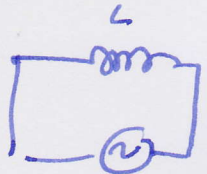


(b)(4) Consider that the **inductor alone** is connected to the power supply, determine the expression of the potential difference across it.

Kirchhoff's loop rule leads to

$$\Delta V_L - L \frac{dI}{dt} = 0 \Rightarrow \Delta V = \omega L I_{\max} \cos \omega t$$

$$\Rightarrow \Delta V = L \omega I_{\max} \sin(\omega t + \pi/2) \Rightarrow \begin{cases} \Delta V_{L, \max} = L \omega I_{\max} \\ \phi = +\pi/2 \end{cases}$$



(c)(6) When both  $R$  and  $L$  are put in series, determine the impedance of the circuit.

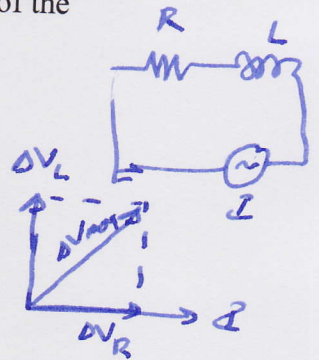
We use the phasor diagram for

$$\Delta V = \Delta V_L + \Delta V_R$$

$$\text{or } \Delta V_H = \sqrt{\Delta V_{L, \max}^2 + \Delta V_{R, \max}^2}$$

$$\Delta V_H = \sqrt{R^2 + L^2 \omega^2} \cdot I_{\max}$$

$$\Rightarrow \boxed{Z = \sqrt{R^2 + L^2 \omega^2}}$$



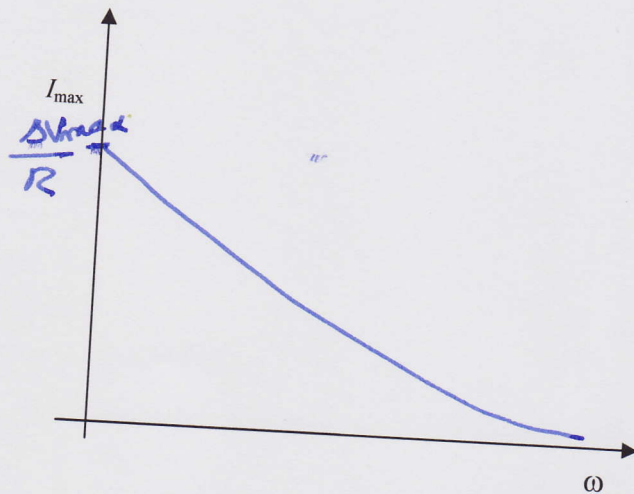
(d)(5) Determine the expression of the circuit phase.

$$\tan \phi = \frac{\Delta V_{L, \max}}{\Delta V_{R, \max}} = \frac{L \omega I_{\max}}{R I_{\max}} \Rightarrow \boxed{\phi = \tan^{-1} \left( \frac{L \omega}{R} \right)}$$

- (e) (6) Use the axes below to plot the current  $I_{\max}$  vs.  $\omega$  assuming  $\Delta V_{\max}$  is constant and discuss the behavior of this circuit at low frequencies ( $\omega \rightarrow 0$ ) and at high frequencies ( $\omega \rightarrow \infty$ ).

$$\text{as } \omega \rightarrow 0 \Rightarrow Z \Rightarrow R \Rightarrow I_{\max} \rightarrow \frac{\Delta V_{\max}}{R}$$

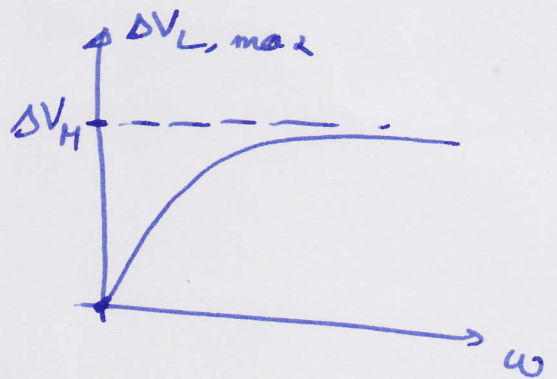
$$\text{as } \omega \rightarrow \infty \Rightarrow Z \rightarrow \infty \Rightarrow I_{\max} \rightarrow 0$$



- Find the relationship between the potential drop across the inductor and the power supply. What might one use this circuit for?
- (f) (5) Determine the voltage across the inductor and plot it as a function of the angular frequency.

$$\Delta V_L = Z_L I \Rightarrow \Delta V_{L, \max} = L\omega I_{\max}$$

$$= \frac{L\omega}{\sqrt{R^2 + L^2\omega^2}} \Delta V_M$$



This circuit can be used as a high-pass filter where oscillations at low frequency are not 'visible' around the inductor.